

# Strange Nucleon Form Factors from $ep$ and $\nu p$ Elastic Scattering

A combined analysis of HAPPEX,  $G^0$ , and BNL E734 data

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# Outline

- Parity-violating electron-nucleon (PVeN) elastic scattering
- Neutrino-proton elastic scattering
- How the use of neutrino-proton elastic scattering data with PVeN data permits the extraction of the strange vector and axial form factors
- The results of a combined analysis of BNL E734  $\nu p$  data with the HAPPEX and  $G^0$  forward PVeN data begin to show the  $Q^2$ -dependence of the strange axial form factor for  $Q^2 = 0.45 \text{ -- } 1.0 \text{ GeV}^2$ .
- **These results greatly restrict the possible configurations of strange quarks in the nucleon.**
- **How these results will aid the determination of the  $Q^2$ -dependence of the strange vector form factors**

# Features of parity-violating forward-scattering $ep$ data

- measures linear combination of form factors of interest
  - axial terms are doubly suppressed
    - $(1 - 4\sin^2\theta_W) \sim 0.075$
    - kinematic factor  $\varepsilon' \sim 0$  at forward angles
  - significant radiative corrections exist, especially in the axial term
- ▣▶ parity-violating data at forward angles are mostly sensitive to the strange electric and magnetic form factors

# Full Expression for the PV $ep$ Asymmetry

For a hydrogen target, the asymmetry as a linear combination of  $G_E^s$ ,  $G_M^s$ ,  $G_A^{CC}$  and  $G_A^s$  is:

$$A^p = A_0^p + A_E^p G_E^s + A_M^p G_M^s + A_{AIV}^p G_A^{CC} + A_A^p G_A^s$$

$$\text{where } A_0^p = -K^p \left\{ \begin{array}{l} (1 - 4\sin^2 \theta_w)(1 + R_V^p) (\epsilon G_E^{p^2} + \tau G_M^{p^2}) \\ - (1 + R_V^n) (\epsilon G_E^p G_E^n + \tau G_M^p G_M^n) \\ - \epsilon' G_M^p (1 - 4\sin^2 \theta_w) [\sqrt{3} R_A^{T=0} G_A^8] \end{array} \right\}$$

$$A_E^p = K^p \left\{ \epsilon G_E^p (1 + R_V^0) \right\}$$

$$A_M^p = K^p \left\{ \tau G_M^p (1 + R_V^0) \right\}$$

$$A_{AIV}^p = K^p \left\{ \epsilon' G_M^p (1 - 4\sin^2 \theta_w) (1 + R_A^{T=1}) \right\}$$

$$A_A^p = K^p \left\{ \epsilon' G_M^p (1 - 4\sin^2 \theta_w) (1 + R_A^0) \right\}$$

$$K^p = \frac{G_F Q^2}{4\pi\sqrt{2}\alpha} \frac{1}{\epsilon G_E^{p^2} + \tau G_M^{p^2}}$$

$$\tau = \frac{Q^2}{4M^2}$$

$$\epsilon = \left[ 1 + 2(1 + \tau) \tan^2(\theta/2) \right]^{-1}$$

$$\epsilon' = \sqrt{(1 - \epsilon^2)\tau(1 + \tau)}$$

Note suppression of axial terms by  $(1 - 4\sin^2 \theta_w)$  and  $\epsilon'$ .

# Things known and unknown in the PV $ep$ Asymmetry

$G_{E,M}^{p,n}$  = Kelly parametrization [PRC 70 (2004) 068202]

with  $G^0$  uncertainties [<http://www.npl.uiuc.edu/exp/G0/Forward>]

$$G_A^{CC} = \frac{g_A}{(1 + Q^2/M_A^2)^2} \quad G_A^8 = \frac{1}{2\sqrt{3}} \frac{(3F - D)}{(1 + Q^2/M_A^2)^2}$$

$M_A = 1.001 \pm 0.020$  GeV [Budd, Bodek and Arrington : hep - ex/0308005 and 0410055]

$g_A = 1.2695 \pm 0.0029$  [Particle Data Group 2005]

$3F - D = 0.585 \pm 0.025$  [Goto *et al.* PRD 62 (2000) 034017]

[use of  $3F - D$  implies use of flavor - SU(3), but  $G_A^8$  is suppressed by  $\varepsilon'$  and  $(1 - 4\sin^2 \theta_w)$ ]

The  $R$ 's are radiative corrections calculated at  $Q^2 = 0$  in the formalism of Zhu et al. [PRD 62 (2000) 033008]. The  $Q^2$  - dependence is unknown,

and so we have assigned a 100% uncertainty to the values.

$$R_V^p = -0.045 \quad R_V^n = -0.012 \quad R_V^0 = -0.012$$

$$R_A^{T=1} = -0.173 \quad R_A^{T=0} = -0.253 \quad R_A^0 = -0.552$$

[from evaluation of Arvieux *et al.*, to be published]

# Features of elastic $\nu p$ data

- measures quadratic combination of form factors of interest
- axial terms are dominant at low  $Q^2$

$$\frac{d\sigma}{dQ^2}(\nu p \rightarrow \nu p) \xrightarrow{Q^2 \rightarrow 0} \frac{G_F^2}{128\pi} \frac{M_p^2}{E_\nu^2} \left[ \left( -G_A^u + G_A^d + G_A^s \right)^2 + \left( 1 - 4 \sin^2 \theta_w \right)^2 \right]$$

- radiative corrections are insignificant

[Marciano and Sirlin, PRD 22 (1980) 2695]

▣► neutrino data are mostly sensitive to the strange axial form factor

# Elastic NC neutrino-proton cross sections

$$\frac{d\sigma}{dQ^2}(\nu p \rightarrow \nu p) = \frac{G_F^2}{2\pi} \frac{Q^2}{E_\nu^2} \left( A \pm BW + CW^2 \right) \quad \begin{array}{l} + \nu \\ - \bar{\nu} \end{array}$$

$$W = 4(E_\nu/M_p - \tau) \quad \tau = Q^2/4M_p^2$$

$$A = \frac{1}{4} \left[ \left( G_A^Z \right)^2 (1 + \tau) - \left( \left( F_1^Z \right)^2 - \tau \left( F_2^Z \right)^2 \right) (1 - \tau) + 4\tau F_1^Z F_2^Z \right]$$

$$B = -\frac{1}{4} G_A^Z \left( F_1^Z + F_2^Z \right)$$

$$C = \frac{1}{64\tau} \left[ \left( G_A^Z \right)^2 + \left( F_1^Z \right)^2 + \tau \left( F_2^Z \right)^2 \right]$$

Dependence on strange form factors is buried in the weak (Z) form factors.

# The BNL E734 Experiment

- performed in mid-1980's
- measured neutrino- and antineutrino-proton elastic scattering
- used wide band neutrino and anti-neutrino beams of  $\langle E_\nu \rangle = 1.25$  GeV
- covered the range  $0.45 < Q^2 < 1.05$  GeV<sup>2</sup>
- large liquid-scintillator target-detector system
- still the only elastic neutrino-proton cross section data available

# E734

Uncertainties shown are total (stat and sys).

## Results

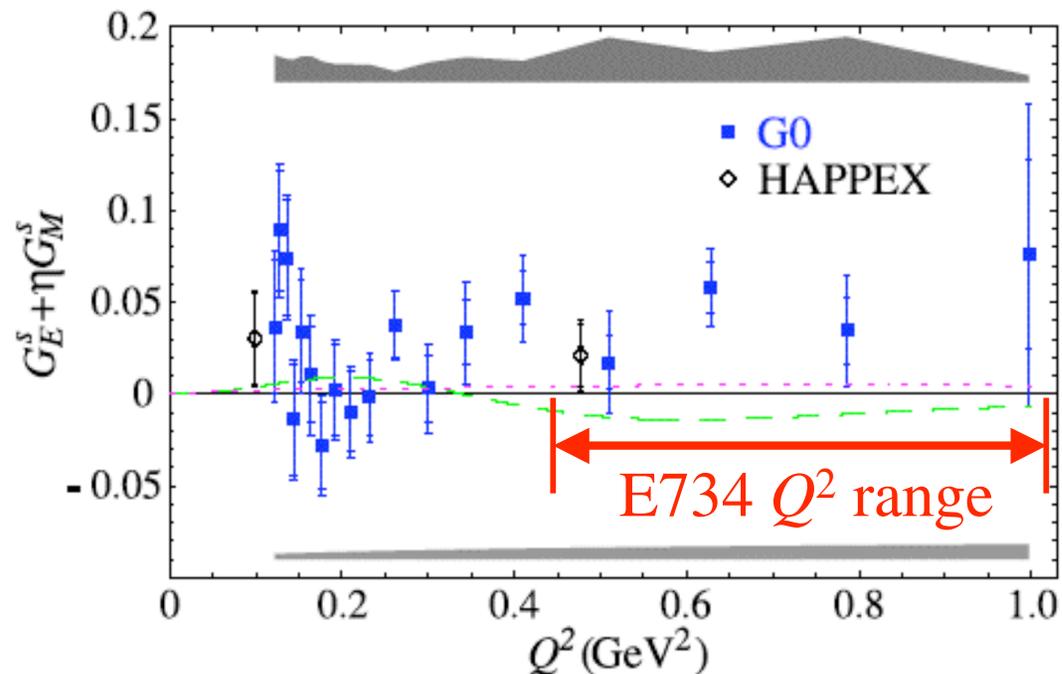
Correlation coefficient arises from systematic errors.

$Q^2$ (GeV) <sup>2</sup>	$d\sigma/dQ^2(\nu p)$ (fm/GeV) <sup>2</sup>	$d\sigma/dQ^2(\bar{\nu} p)$ (fm/GeV) <sup>2</sup>	correlation coefficient
0.45	0.165 ± 0.033	0.0756 ± 0.0164	0.134
0.55	0.109 ± 0.017	0.0426 ± 0.0062	0.256
0.65	0.0803 ± 0.0120	0.0283 ± 0.0037	0.294
0.75	0.0657 ± 0.0098	0.0184 ± 0.0027	0.261
0.85	0.0447 ± 0.0092	0.0129 ± 0.0022	0.163
0.95	0.0294 ± 0.0074	0.0108 ± 0.0022	0.116
1.05	0.0205 ± 0.0062	0.0101 ± 0.0027	0.071

# Combination of the $ep$ and $\nu p$ data sets

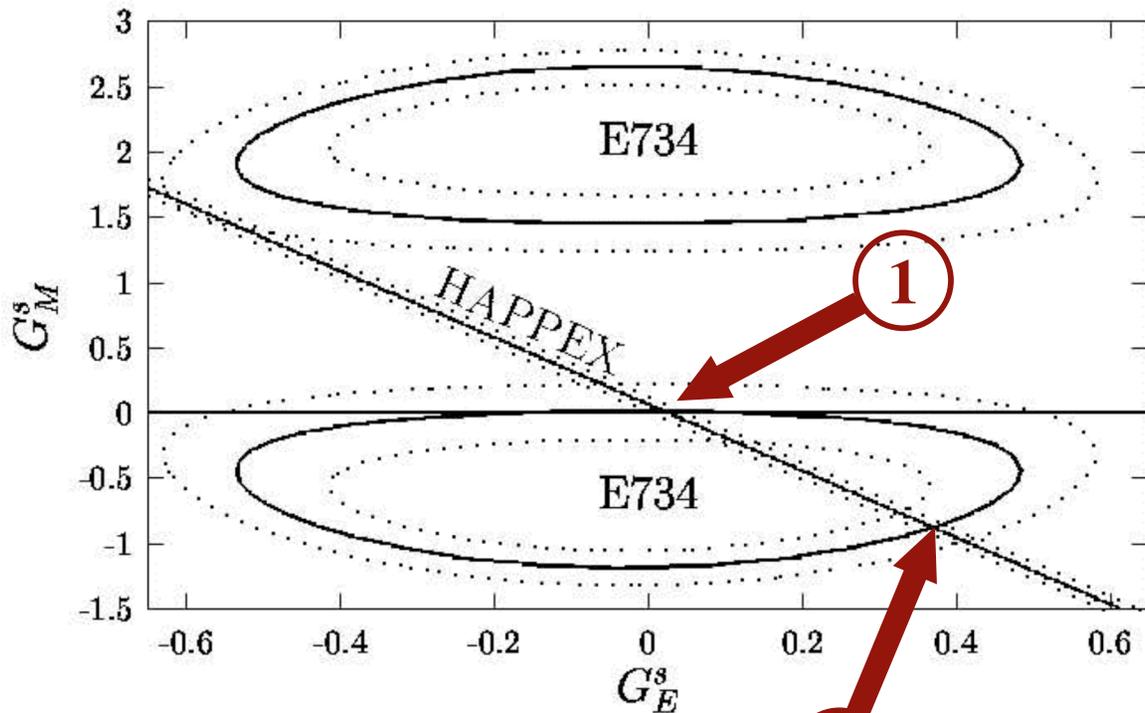
We use PV  $ep$  data in the same range of  $Q^2$  as the E734 experiment.

- The original HAPPEX measurement:  $Q^2 = 0.477 \text{ GeV}^2$   
[PLB 509 (2001) 211 and PRC 69 (2004) 065501]
- The recent  $G^0$  data covering the range  $0.1 < Q^2 < 1.0 \text{ GeV}^2$   
[PRL 95 (2005) 092001]



# Combination of the $ep$ and $\nu p$ data sets

Since the neutrino data are quadratic in the form factors, then there will be in general two solutions when these data sets are combined.



**Here are the advertised intersections between electron and neutrino data.**

$$Q^2 = 0.5 \text{ GeV}^2$$

Fortunately, the two solutions are very distinct from each other, and other available data can select the correct physical solution.

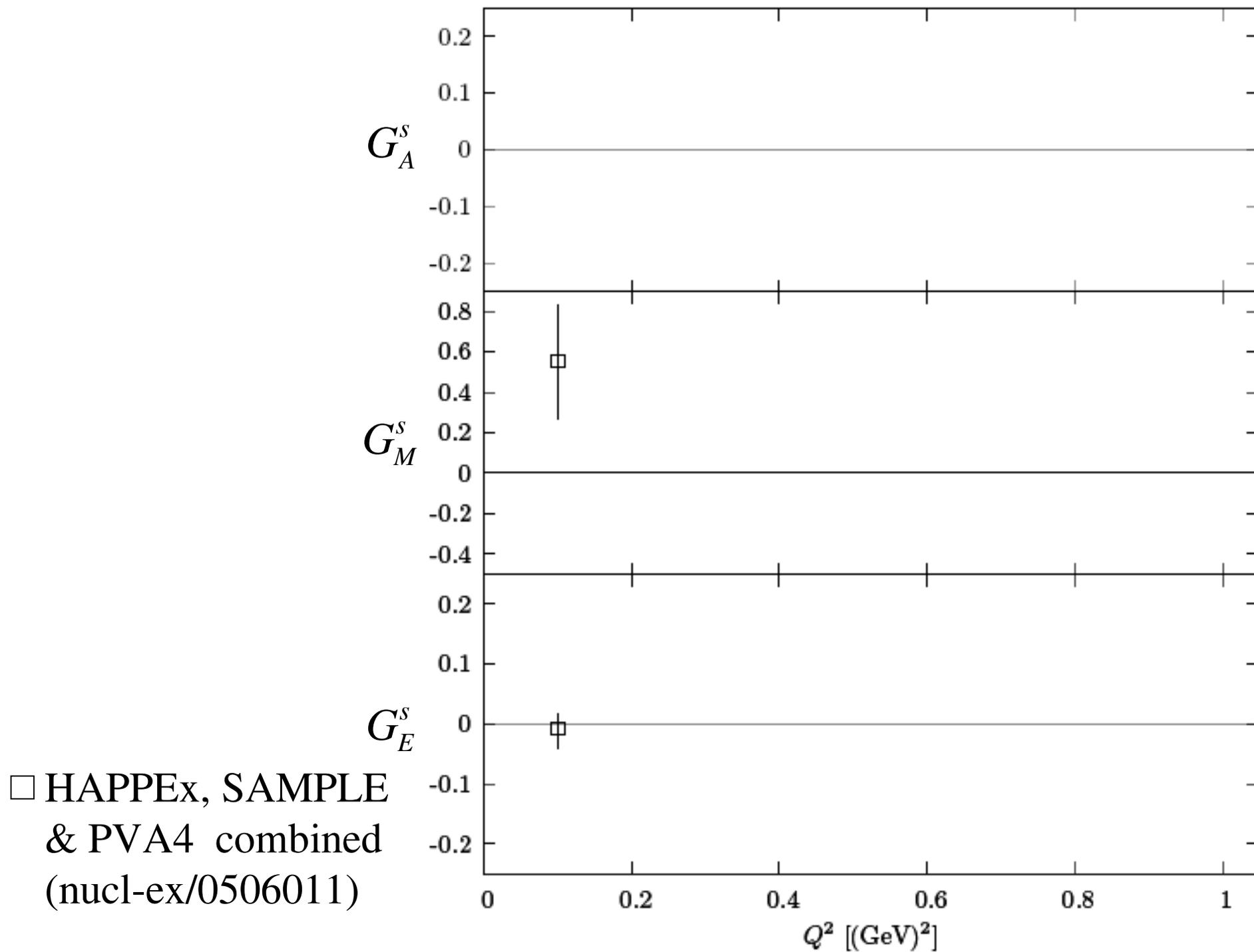
## General Features of the two Solutions

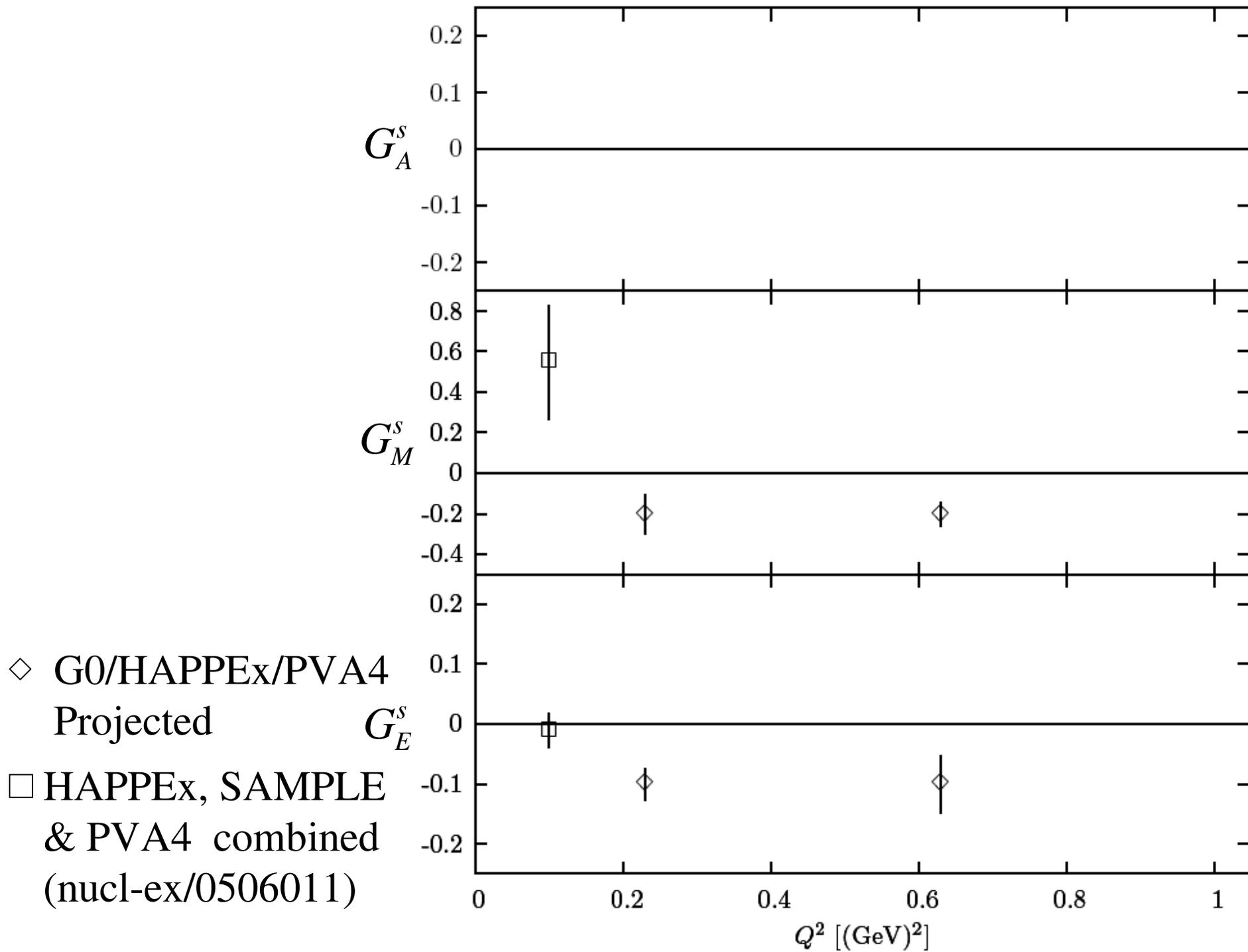
	Solution 1	Solution 2
$G_E^s$	Consistent with zero (with large uncertainty)	Large and positive
$G_M^s$	Consistent with zero (with large uncertainty)	Large and negative
$G_A^s$	Small and negative	Large and positive

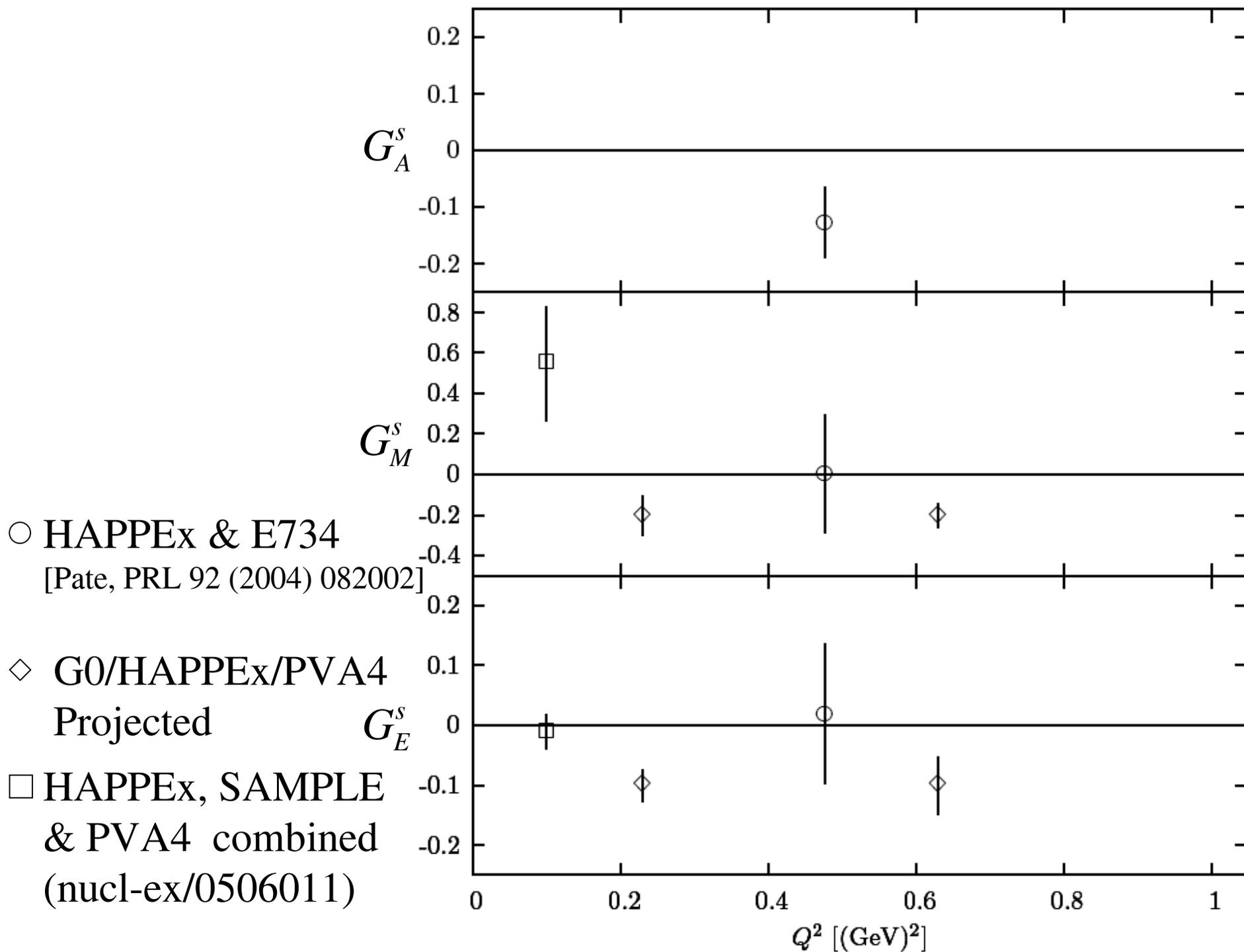
There are three strong reasons to prefer Solution 1:

- $G_A^s$  in Solution 2 is inconsistent with DIS estimates for  $\Delta s$
- $G_M^s$  in Solution 2 is inconsistent with the combined SAMPLE/PVA4/HAPPEX result of  $G_M^s = \sim +0.6$  at  $Q^2 = 0.1 \text{ GeV}^2$
- $G_E^s$  in Solution 2 is inconsistent with the idea that  $G_E^s$  should be small, and conflicts with expectation from recent  $G^0$  data that  $G_E^s$  may be negative near  $Q^2 = 0.3 \text{ GeV}^2$

I only present Solution 1 in what follows.

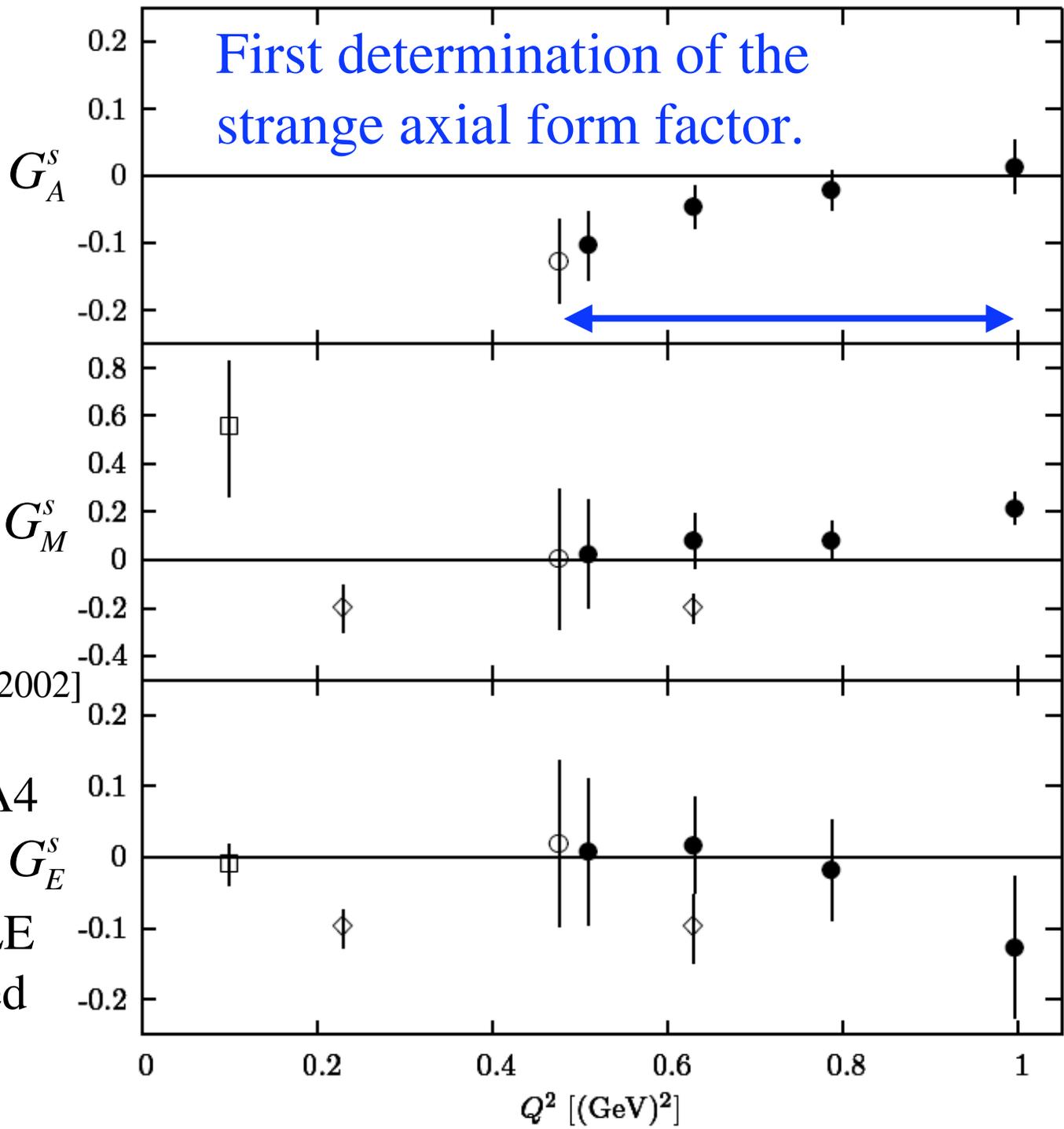


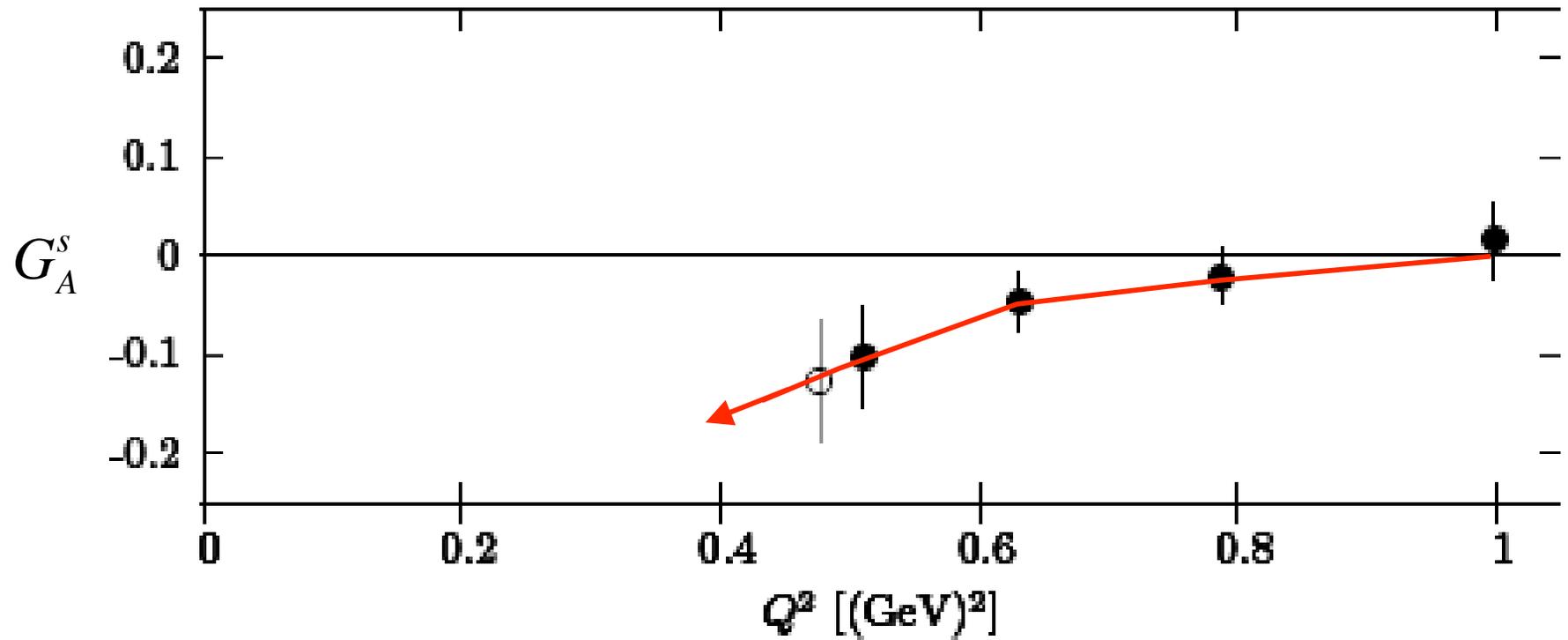




First determination of the strange axial form factor.

- G0 & E734  
[to be published]
- HAPPE<sub>x</sub> & E734  
[Pate, PRL 92 (2004) 082002]
- ◇ G0/HAPPE<sub>x</sub>/PVA4  
Projected
- HAPPE<sub>x</sub>, SAMPLE & PVA4 combined  
(nucl-ex/0506011)

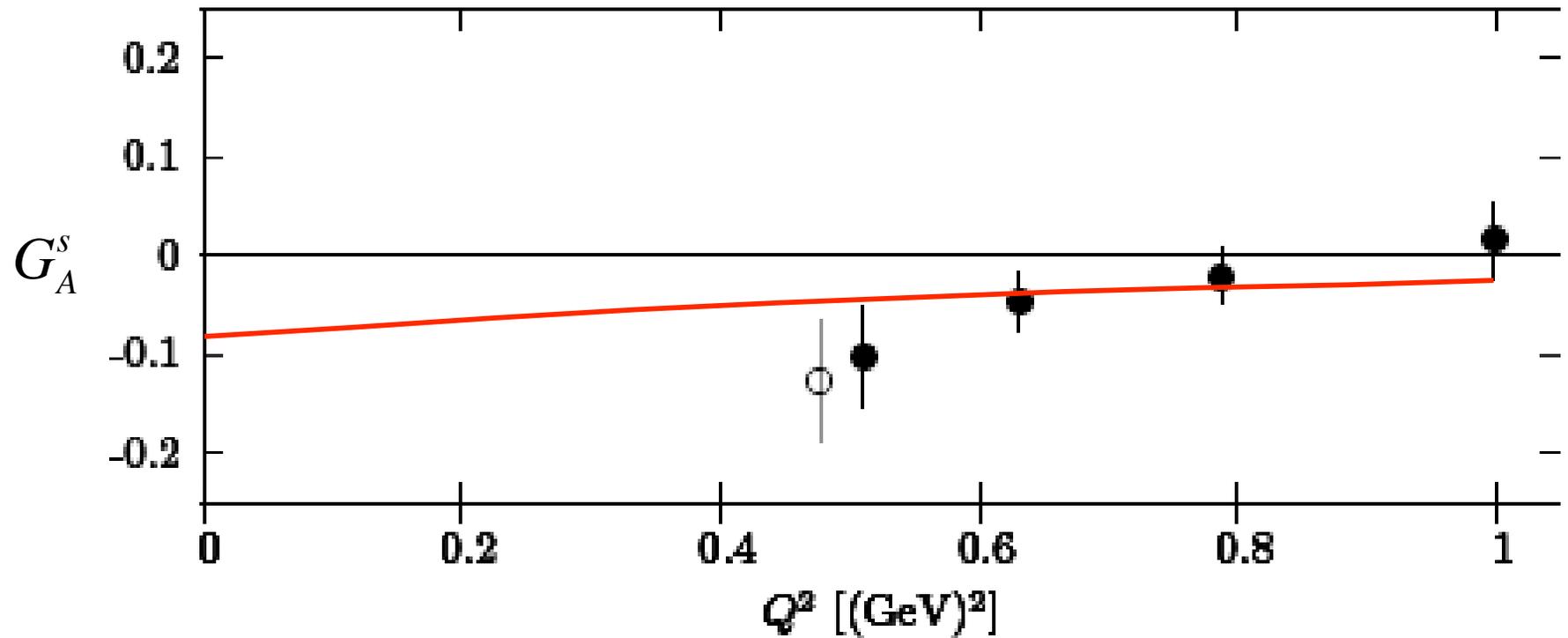




- G0 & E734  
[to be published]

- HAPPEX & E734  
[Pate, PRL 92 (2004) 082002]

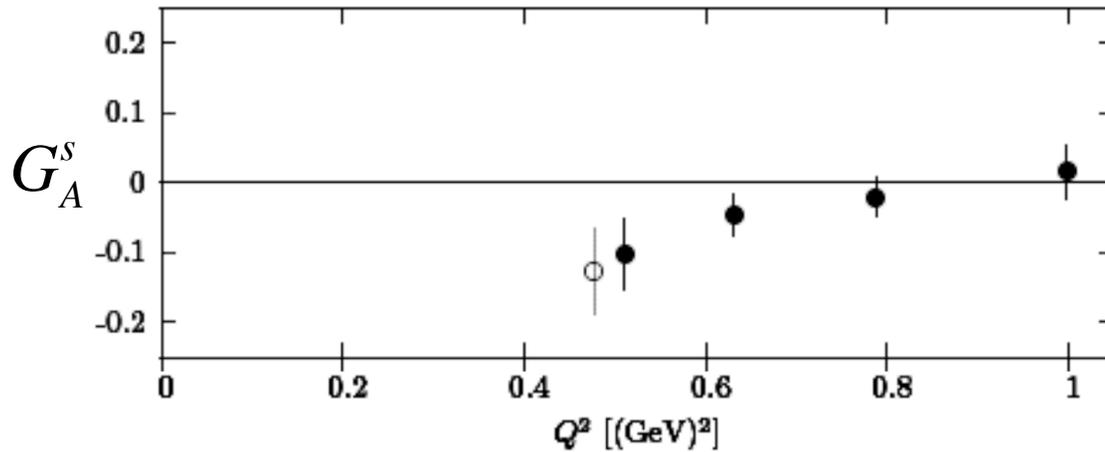
$Q^2$ -dependence suggests  $\Delta s < 0$  !



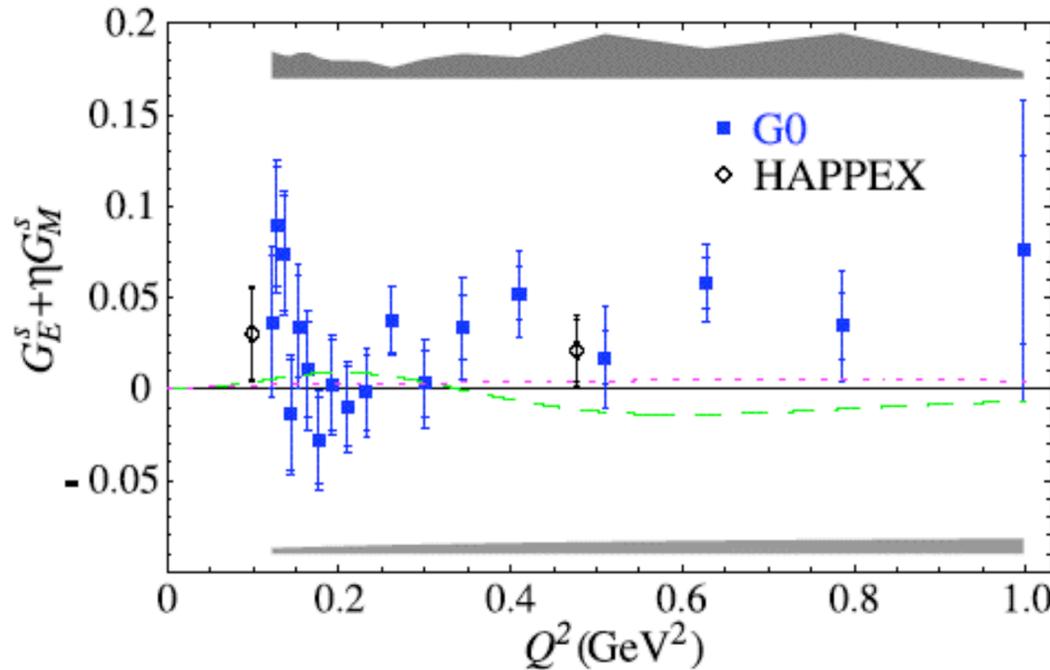
● G0 & E734  
[to be published]

○ HAPPE<sub>x</sub> & E734  
[Pate, PRL 92 (2004) 082002]

Recent calculation by Silva, Kim, Urbano, and Goeke (hep-ph/0509281 and Phys. Rev. D 72 (2005) 094011) based on chiral quark-soliton model is in rough agreement with the data.



These results on  $G_A^s$ ...



combined with world data  
on  $G_E^s + \eta G_M^s$ ...

determine a unique  $uuds\bar{s}$   
configuration, in which  
the  $uuds$  system is  
radially excited and the  $\bar{s}$   
is in the ground state.

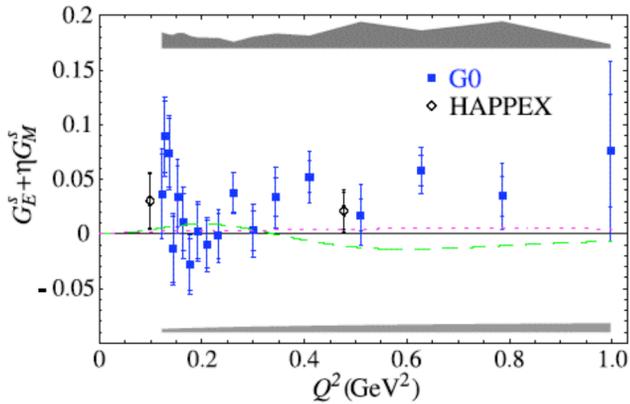
An, Riska and Zou, hep-ph/0511223; Riska and Zou, nucl-th/0512102.

# Strange Vector Form Factors: Using $ep$ and $\nu p$ data

The international program of PV  $ep$  measurements will completely resolve the strange vector form factors at only three  $Q^2$  points: 0.1, 0.23 and 0.63  $\text{GeV}^2$ .

At many other points in the range  $0.038 < Q^2 < 1.0 \text{ GeV}^2$ , we have PV physics asymmetries that represent linear combinations of the vector, axial and anapole form factors. These can be used to constrain fits that seek to understand the  $Q^2$ -dependence of these form factors.

As has just been demonstrated, a combination of the forward-scattering PV  $ep$  data with elastic  $\nu p$  data provides several more points where the vector form factors are resolved, and with reasonable error bars: 0.63, 0.79, and 0.99  $\text{GeV}^2$ . **These can already provide powerful constraints on global fits.**



Global fits will use these five resolved points,  
and also the linear combination constraints.



● G0 & E734  
[to be published]

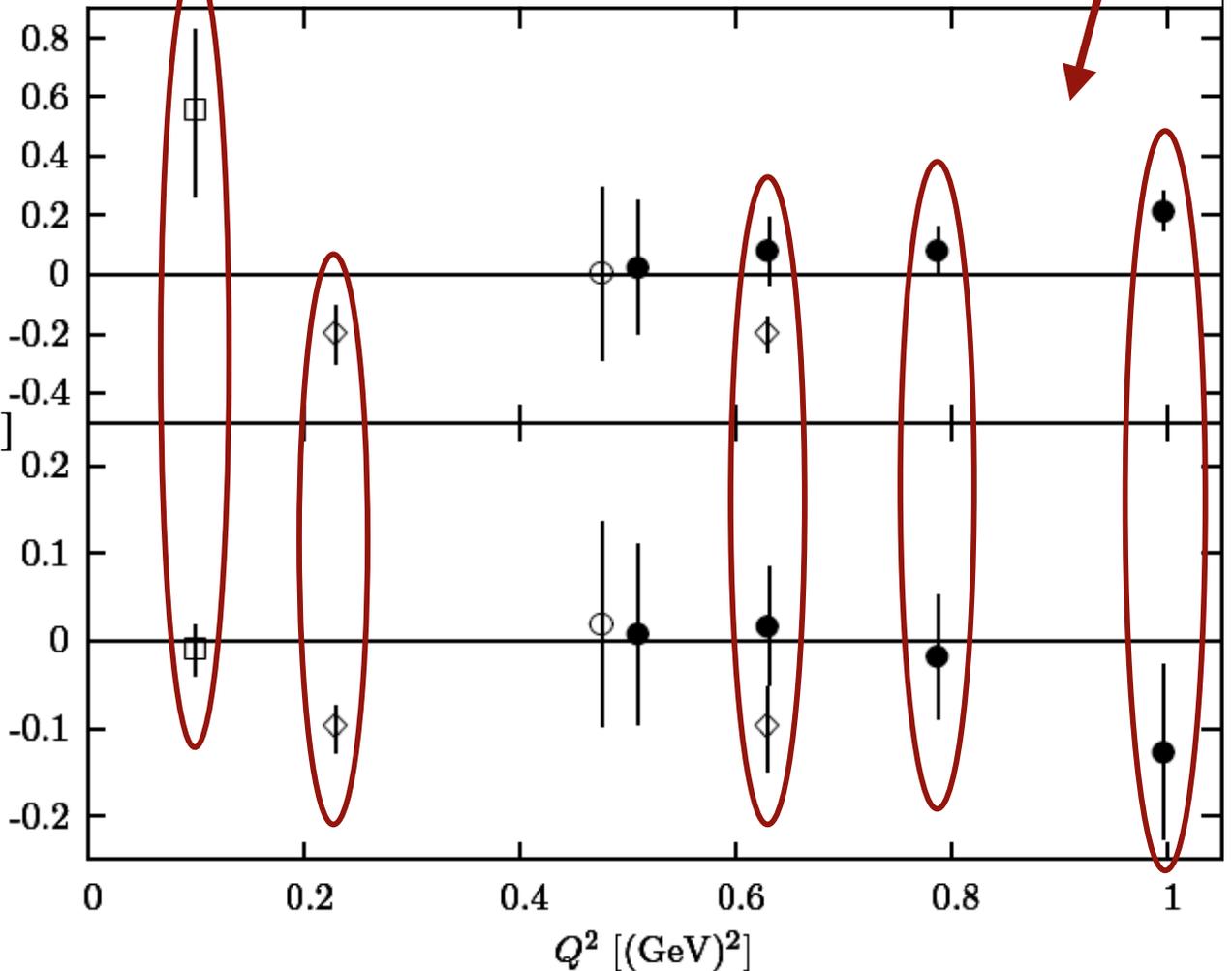
○ HAPPEX & E734  
[Pate, PRL 92 (2004) 082002]

◇ G0/HAPPEX/PVA4  
Projected

□ HAPPEX, SAMPLE  
& PVA4 combined  
(nucl-ex/0506011)

$G_M^s$

$G_E^s$



But better  $\nu p$  data are needed!

The E734 data have insufficient precision and too narrow a  $Q^2$  range to achieve the full potential of this physics program. Better neutrino data is needed, with smaller uncertainties and points nearer  $Q^2 = 0$ , to fulfill the potential of this analysis method.

A detailed understanding of the  $Q^2$ -dependence of these form factors will not be possible until a more dense set of resolved data points are available. FINE SSE can provide these additional data points.

△ **FINeSSE (& G0)**  
 [exp. proposal:  
 no nuclear initial or  
 final state effects  
 included in errors]

● **G0 & E734**  
 [to be published]

○ **HAPPE<sub>x</sub> & E734**  
 [Pate, PRL 92 (2004) 082002]

◇ **G0/HAPPE<sub>x</sub>/PVA4**  
**Projected**

□ **HAPPE<sub>x</sub>, SAMPLE  
 & PVA4 combined**  
 (nucl-ex/0506011)

